

Shape optimization of a layer-by-layer constraint functional for additive manufacturing processes

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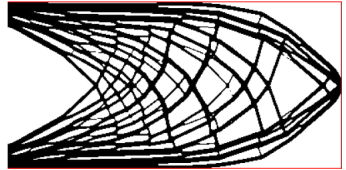
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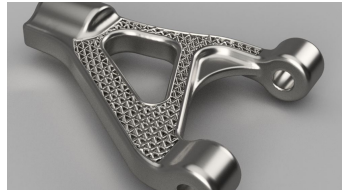
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Foreword: shape optimization in the industrial context

- **Shape and topology optimization** techniques have aroused a tremendous enthusiasm within the engineering and industrial communities.
- One drawback of these methods is that the optimized designs are often too complicated to be constructed by traditional methods such as **milling** or **casting**.
- The recent headway made by **additive manufacturing** methods allow to assemble structures with a high degree of complexity.
- These techniques impose new constraints on the manufactured components.



Typical 'truss' designs resulting from shape and topology optimization processes.

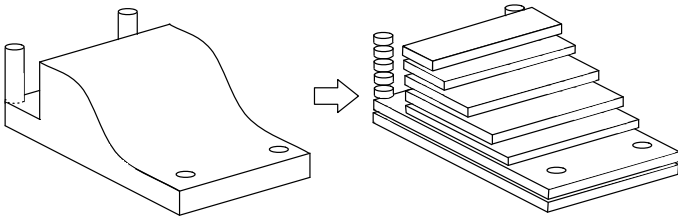


Part produced with an additive manufacturing method (from <http://www.autodesk.com/>).

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Additive manufacturing in a nutshell (I)

- All the **additive manufacturing processes** begin with a **slicing** stage: the input shape is decomposed into a series of **horizontal layers**.
- These $2d$ layers are built one on top of the other according to the selected technology.

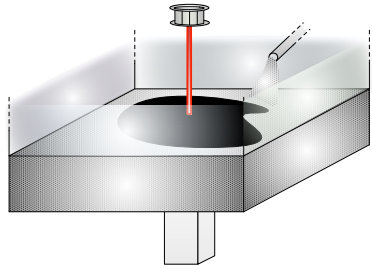
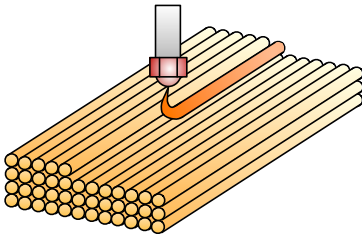


Sketch of the **slicing procedure**, initiating any additive manufacturing process.

Additive manufacturing in a nutshell (II)

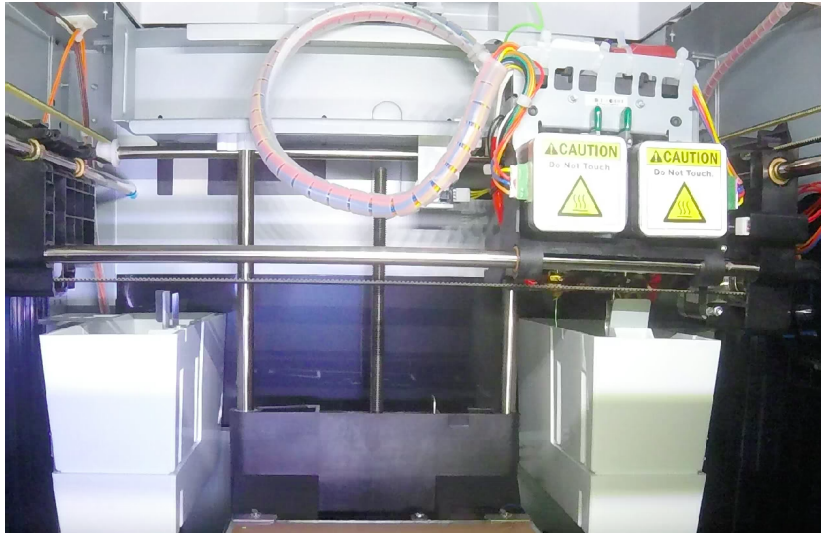
Two popular additive manufacturing technologies are:

- **Material extrusion methods** (e.g. FDM): they act by **direct deposition** of rasters of a molten filament. Such methods are often used to process plastic (ABS).
- **Powder bed fusion methods** (e.g. EBM, SLS), which are generally used to process metals. Each 2d layer is assembled by spreading metallic powder within the build chamber, then binding the grains together with a laser.



*Sketch of the (left) **FDM** and (right) **EBM** additive manufacturing processes.*

Additive manufacturing techniques in a nutshell (III)



An FDM machine tool in action.

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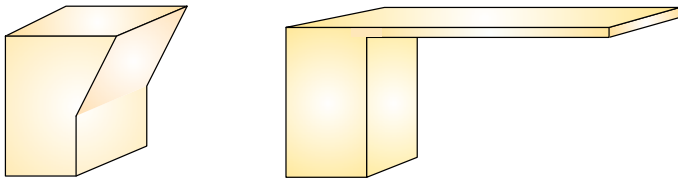
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The overhang issue (I)

All additive manufacturing technologies experience trouble in the assembly of shapes showing large **overhangs**, i.e. regions hanging over void.

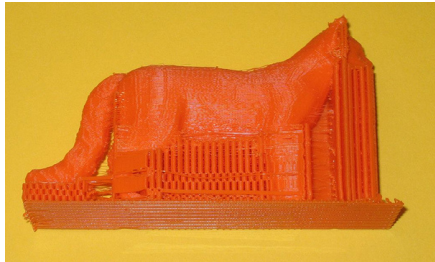
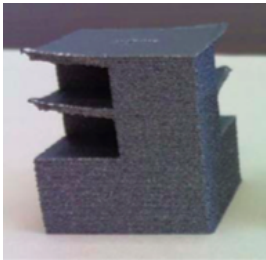
- in the case of FDM processes, this amounts to assembling over void.
- In the case of powder-bed methods, the rapid melting then solidification of the powder induces **residual stress**, especially in regions unanchored to the lower structure. This may cause **warpage** of such parts upon cooling.



(Left) short overhang; support from the lower structure is sufficient to guarantee manufacturability; (right) large overhang.

The overhang issue (II)

- The most common strategy to deal with overhangs is to erect a sacrificial **scaffold structure** alongside the construction of the shape [DuHeLe].
- This scaffold structure has to be removed as a post-processing, which is costly and cumbersome.



(Left) Warpage caused by residual constraints in an EBM assembly (from [CheLuChou]),
(right) scaffold structure in the construction of a part (from <https://hyrulefoundry.wordpress.com/>).

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Shape optimization of linear elastic shapes (I)

In the context of its **final use**, a **shape** is a bounded domain $\Omega \subset \mathbb{R}^d$, which is

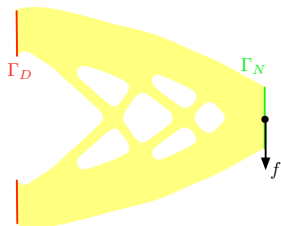
- **fixed** on a part Γ_D of its boundary,
- submitted to **surface loads** f , applied on $\Gamma_N \subset \partial\Omega$, $\Gamma_D \cap \Gamma_N = \emptyset$.

The displacement vector field $u_\Omega : \Omega \rightarrow \mathbb{R}^d$ is governed by the **linear elasticity system**:

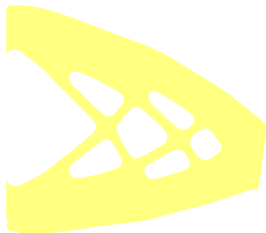
$$\begin{cases} -\operatorname{div}(Ae(u_\Omega)) &= 0 & \text{in } \Omega \\ u_\Omega &= 0 & \text{on } \Gamma_D \\ Ae(u_\Omega)n &= f & \text{on } \Gamma_N \\ Ae(u_\Omega)n &= 0 & \text{on } \Gamma \end{cases},$$

where $e(u) = \frac{1}{2}(\nabla u^T + \nabla u)$ is the **strain tensor**, and A is the **Hooke's law** of the material:

$$\forall e \in \mathcal{S}_d(\mathbb{R}), \quad Ae = 2\mu e + \lambda \operatorname{tr}(e)I.$$



A 'Cantilever'



The deformed cantilever

The shape optimization problem

The shape optimization problem of interest reads:

$$\min_{\mathcal{U}_{\text{ad}}} J(\Omega), \text{ s.t. } P(\Omega) \leq \alpha,$$

in which

- \mathcal{U}_{ad} is a set of (smooth) admissible shapes,
- The objective function $J(\Omega)$ is the **structural compliance** of shapes:

$$J(\Omega) = \int_{\Omega} A e(u_{\Omega}) : e(u_{\Omega}) \, dx = \int_{\Gamma_N} f \cdot u_{\Omega} \, ds,$$

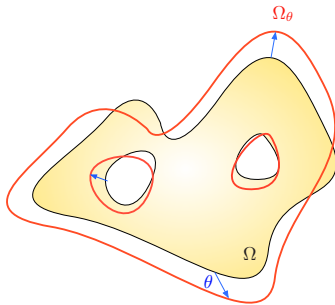
- The constraint $P(\Omega)$ enforces the constructibility by additive manufacturing processes,
- Other constraints may be added to the problem, e.g. on the volume $\text{Vol}(\Omega)$ of shapes.

Differentiation with respect to the domain: Hadamard's method

Hadamard's boundary variation method describes variations of a reference, Lipschitz domain Ω of the form:

$$\Omega \rightarrow \Omega_\theta := (\text{Id} + \theta)(\Omega),$$

for 'small' $\theta \in W^{1,\infty}(\mathbb{R}^d, \mathbb{R}^d)$.



Definition 1.

Given a smooth domain Ω , a function $J(\Omega)$ of the domain is **shape differentiable** at Ω if the function

$$W^{1,\infty}(\mathbb{R}^d, \mathbb{R}^d) \ni \theta \mapsto J(\Omega_\theta)$$

is Fréchet-differentiable at 0, i.e. the following expansion holds around 0:

$$J(\Omega_\theta) = J(\Omega) + J'(\Omega)(\theta) + o(\|\theta\|_{W^{1,\infty}(\mathbb{R}^d, \mathbb{R}^d)}).$$

Differentiation with respect to the domain: Hadamard's method

Techniques from optimal control theory make it possible to calculate shape derivatives; in the case of 'many' functionals of the domain $J(\Omega)$, the shape derivative has the particular **structure**:

$$J'(\Omega)(\theta) = \int_{\Gamma} v_{\Omega} \theta \cdot n \, ds,$$

where v_{Ω} is a scalar field depending on u_{Ω} , and possibly on an **adjoint state** p_{Ω} .

Example: If the objective function

$$J(\Omega) = \int_{\Gamma_N} f \cdot u_{\Omega} \, ds$$

is the **compliance**, $v_{\Omega} = -Ae(u_{\Omega}) : e(u_{\Omega})$ is the (negative) elastic energy density.

The generic algorithm

This shape gradient provides a natural **descent direction** for $J(\Omega)$: *for instance*, defining θ as

$$\theta = -v_\Omega n$$

yields, for $t > 0$ sufficiently small (*to be found numerically*):

$$J((\text{Id} + t\theta)(\Omega)) = J(\Omega) - t \int_\Gamma v_\Omega^2 ds + o(t) < J(\Omega)$$

Gradient algorithm: For $n = 0, \dots$ until convergence,

1. Compute the solution u_{Ω^n} (and p_{Ω^n}) of the elasticity system on Ω^n .
2. Compute the shape gradient $J'(\Omega^n)$ thanks to the previous formula, and infer a descent direction θ^n for the cost functional.
3. **Advect** the shape Ω^n according to θ^n , so as to get $\Omega^{n+1} := (\text{Id} + \theta^n)(\Omega^n)$.

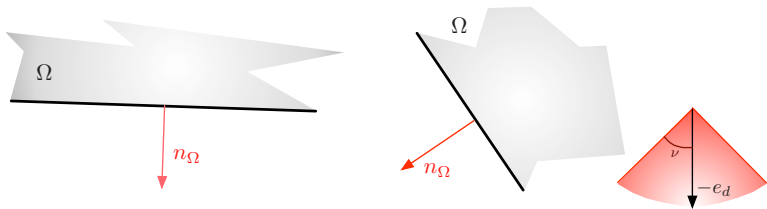
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The 'naive', geometric attempt (I)

- Most approaches in the literature rely on the **angle** between $\partial\Omega$ and the (vertical) build direction to detect and penalize overhangs.
- An intuitive approach relies on **anisotropic perimeter** functionals of the form:

$$P_g(\Omega) = \int_{\partial\Omega} \varphi(n_\Omega) ds, \text{ where } \varphi : \mathbb{R}^d \rightarrow \mathbb{R} \text{ is given.}$$

Example The choice $\varphi_a(n) := (n \cdot e_d + \cos \nu)_-^2$, where $(s)_- := \min(s, 0)$, penalizes regions of $\partial\Omega$ where the angle $n \cdot (-e_d)$ is smaller than a **threshold** ν .



Parts of $\partial\Omega$ (left) violating and (right) satisfying the angle-based criterion.

The 'naive', geometric attempt (II)

Proposition 1.

The functional $P_g(\Omega)$ is *shape differentiable* at any admissible shape $\Omega \in \mathcal{U}_{\text{ad}}$, and its shape derivative reads:

$$P'_g(\Omega)(\theta) = \int_{\Gamma} \kappa \varphi(n) \theta \cdot n \, ds - \int_{\Gamma} \nabla(\varphi(n)) \cdot \nabla_{\partial\Omega}(\theta \cdot n) \, ds,$$

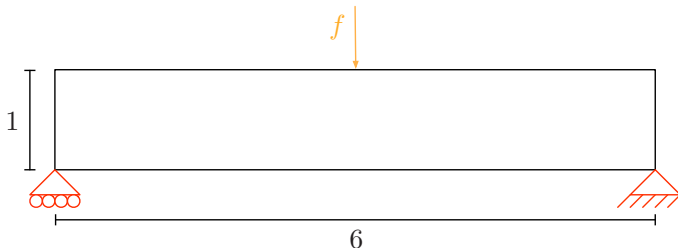
where $\nabla_{\partial\Omega}\psi := \nabla\psi - (\nabla\psi \cdot n)n$ is the *tangential gradient* of a smooth enough function $\psi : \partial\Omega \rightarrow \mathbb{R}$.

- Unfortunately, this approach gives unsatisfactory results.
- We propose instead a general idea for modeling overhang constraints, which appeals to their **mechanical origin**.

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Geometric constraints; the 'dripping effect' (I)

We consider the two-dimensional **MBB Beam** example.



Setting of the two-dimensional MBB beam example.

We first solve the **compliance minimization** problem:

$$\begin{aligned} \min_{\Omega} \quad & J(\Omega), \\ \text{s.t.} \quad & \text{Vol}(\Omega) \leq \alpha_v \text{Vol}(D). \end{aligned}$$

Geometric constraints; the 'dripping effect' (II)



(Top) initial shape Ω_0 and (bottom) optimized shape Ω^* for compliance minimization in the two-dimensional MBB Beam example.

The optimized shape Ω^* presents **large nearly horizontal bars** which are very important for the structural performance.

Geometric constraints; the 'dripping effect' (III)

To help in removing these overhangs, we rather solve the problem:

$$\begin{aligned} \min_{\Omega} \quad & (1 - \alpha_g) \frac{J(\Omega)}{J(\Omega^*)} + \alpha_g \frac{P_g(\Omega)}{P_g(\Omega^*)}, \\ \text{s.t.} \quad & \text{Vol}(\Omega) \leq \alpha_v \text{Vol}(D). \end{aligned}$$



Optimized shape using $\alpha_g = 0.5$.

The shape develops an **oscillatory boundary** so that:

- The angle requirement is (approximately) satisfied,
- The structural performance is not too much altered: the large bars connecting loads to anchor points have not disappeared.

Geometric constraints; the 'dripping effect' (IV)

Adding a **perimeter constraint** to the problem is a tentative remedy, but does not prove sufficient to deal with this issue.



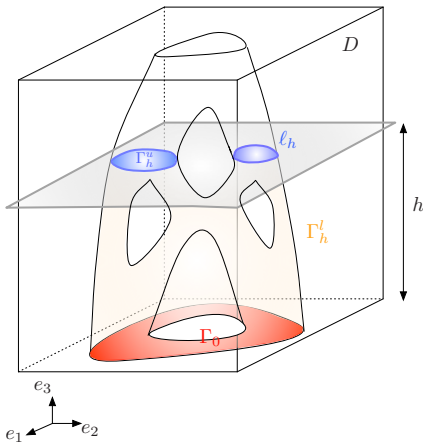
Optimized shape using angle and perimeter penalization.

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Definition of the mechanical constraint (I)

The **mechanical** constraint $P(\Omega)$ relies on the physical behavior of the shape at each stage of its construction.

- Ω is enclosed in the **build chamber** $D = S \times (0, H)$, where $S \subset \mathbb{R}^{d-1}$,
- $\Omega_h := \{x = (x_1, \dots, x_d) \in \Omega, x_d < h\}$ is the **intermediate shape** at height h .
- The boundary $\partial\Omega_h$ is decomposed as $\partial\Omega_h = \Gamma_0 \cup \Gamma_h^u \cup \Gamma_h^l$, where
 - $\Gamma_0 = \{x \in \partial\Omega_h, x_d = 0\}$ is the **contact region** between Ω and the build table,
 - $\Gamma_h^u = \{x \in \partial\Omega_h, x_d = h\}$ is the **upper side** of Ω_h ,
 - $\Gamma_h^l = \partial\Omega_h \setminus (\overline{\Gamma_0} \cup \overline{\Gamma_h^u})$ is the **lateral surface**.



Definition of the mechanical constraint (II)

- Each intermediate shape Ω_h is only subjected to **gravity effects** $g \in H^1(\mathbb{R}^d)^d$. The elastic displacement of Ω_h satisfies:

$$\begin{cases} -\operatorname{div}(Ae(u_{\Omega_h}^c)) = g & \text{in } \Omega_h, \\ u_{\Omega_h}^c = 0 & \text{on } \Gamma_0, \\ Ae(u_{\Omega_h}^c)n = 0 & \text{on } \Gamma_h^l \cup \Gamma_h^u. \end{cases}$$

- The **self-weight** of each intermediate shape Ω_h is:

$$c_{\Omega_h} := \int_{\Omega_h} Ae(u_{\Omega_h}^c) : e(u_{\Omega_h}^c) \, dx = \int_{\Omega_h} g \cdot u_{\Omega_h}^c \, dx.$$

- The **(self-weight) manufacturing compliance** of a **final shape** Ω aggregates the self weights of all its **intermediate shapes**:

$$P_{\text{sw}}(\Omega) = \int_0^H j(c_{\Omega_h}) \, dh,$$

where $j : \mathbb{R} \rightarrow \mathbb{R}$ is a smooth function.

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Shape derivative of the manufacturing compliance (I)

- We consider a fixed shape $\Omega \in \mathcal{U}_{\text{ad}}$.
- Perturbations θ are confined to a class X^k of vector fields of class \mathcal{C}^k , which **identically vanish** near the 'flat regions' of $\partial\Omega$.

Theorem 2.

The manufacturing compliance $P_{\text{sw}}(\Omega)$ is shape differentiable at Ω , in the sense that the mapping $\theta \mapsto P_{\text{sw}}(\Omega_\theta)$, from X^k into \mathbb{R} is differentiable for $k \geq 1$. Its derivative is:

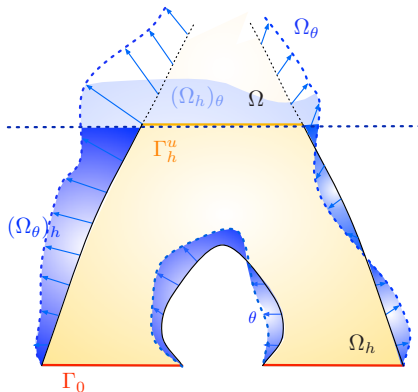
$$\forall \theta \in X^k, \quad P'_{\text{sw}}(\Omega)(\theta) = \int_{\partial\Omega \setminus \overline{\Gamma_0}} \mathcal{D}_\Omega \theta \cdot n \, ds,$$

where the integrand factor \mathcal{D}_Ω is defined, for a.e. $x \in \partial\Omega \setminus \overline{\Gamma_0}$, by:

$$\mathcal{D}_\Omega(x) = \int_{x_d}^H j'(c_{\Omega_h}) (2g \cdot u_{\Omega_h}^c - Ae(u_{\Omega_h}^c) : e(u_{\Omega_h}^c)) (x) \, dh.$$

Shape derivative of the mechanical constraint (II)

Sketch of the proof: The proof is not completely standard since cuts $(\Omega_\theta)_h$ of variations of Ω are not 'classical' variations of shapes.



In general, cuts $(\Omega_\theta)_h$ cannot be described as variations (e.g. $(\Omega_h)_\theta$) of Ω_h .

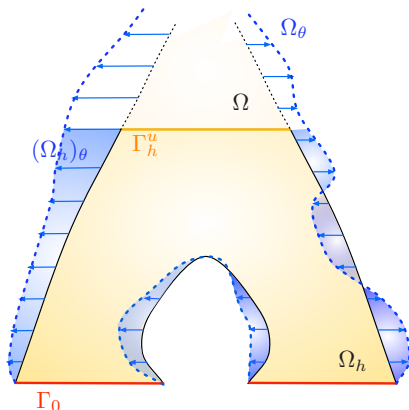
Step I: The case of horizontal perturbations

Step 1 The Theorem holds in the particular case of **horizontal perturbations**

$$\theta \in X_H^k, \quad X_H^k := \left\{ \theta \in X^k, \theta \cdot e_d = 0 \right\}.$$

This part boils down to classical variations of shape, since [in this case](#):

$$(\Omega_\theta)_h = (\Omega_h)_\theta.$$

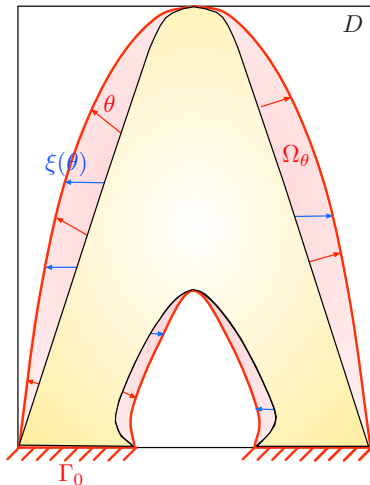


Step II: From 'general' to horizontal perturbations

Step 2 We prove that every variation Ω_θ of Ω with $\theta \in X^{k+1}$ can be described under the form $\Omega_{\xi(\theta)}$, where $\xi(\theta) \in X_H^k$ is horizontal. We also calculate the derivative of the mapping $\theta \mapsto \xi(\theta)$.

This is an application of the [implicit function theorem](#), upon rewriting an equality such as $\Omega_\theta = \Omega_\xi$ under the form

$$\mathcal{G}(\theta, \xi) = 0.$$



Step III: Use of the chain rule

Step III The differentiability of

$$\theta \mapsto P_{\text{sw}}(\Omega_\theta) = P_{\text{sw}}(\Omega_{\xi(\theta)})$$

follows from the chain rule.



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Other models

Other models may be used for the physical behavior of **intermediate shapes** Ω_h . For instance,

- The definition of $u_{\Omega_h}^c$ could be replaced by:

$$\left\{ \begin{array}{ll} -\operatorname{div}(Ae(u_{\Omega_h}^a)) = g_h & \text{in } \Omega_h, \\ u_{\Omega_h}^a = 0 & \text{on } \Gamma_0, \\ Ae(u_{\Omega_h}^a)n = 0 & \text{on } \Gamma_h^l, \\ Ae(u_{\Omega_h}^a)n = 0 & \text{on } \Gamma_h^u, \end{array} \right. \quad \text{where } g_h(x) = \begin{cases} g & \text{if } x_d \in (h - \delta, h), \\ 0 & \text{otherwise,} \end{cases}$$

is an **artificial force** acting on the upper side of Ω_h . As we shall see, this formulation is better at penalizing perfectly horizontal parts hanging over void.

- The mechanical constraint $P(\Omega)$ could involve the solutions v_{Ω_h} to a **thermal cooling problem** posed on Ω_h , to model e.g. **residual stresses** in the final shape Ω ; see [AlJak].

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Numerical evaluation of $P_{\text{sw}}(\Omega)$ and $P'_{\text{sw}}(\Omega)(\theta)$

- The expressions for $P_{\text{sw}}(\Omega)$ and its derivative $P'_{\text{sw}}(\Omega)(\theta)$

$$P_{\text{sw}}(\Omega) = \int_0^H j(c_{\Omega_h}) dh, \text{ and } P'_{\text{sw}}(\Omega)(\theta) = \int_{\partial\Omega \setminus \overline{\Gamma}_0} \mathcal{D}_\Omega \theta \cdot n ds,$$

where

$$\mathcal{D}_\Omega(x) = \int_{x_d}^H j'(c_{\Omega_h}) (2g \cdot u_{\Omega_h}^\varepsilon - Ae(u_{\Omega_h}^\varepsilon) : e(u_{\Omega_h}^\varepsilon)) (x) dh.$$

involve a **continuum** of shapes Ω_h via the self-weights c_{Ω_h} and the elastic displacements $u_{\Omega_h}^\varepsilon$,

- Hence the need for a **suitable discretization** of $h \mapsto c_{\Omega_h}$ and $h \mapsto u_{\Omega_h}^\varepsilon$.
- We present:
 - 0th-order** approximations P_N^0 and \mathcal{D}_N^0 of $P_{\text{sw}}(\Omega)$ and \mathcal{D}_Ω ,
 - First-order** approximations P_N^1 and \mathcal{D}_N^1 based on an **interpolation procedure**.

The 'naive', 0th order method

- The height interval $(0, H)$ is discretized with a sequence:

$$0 < h_1 < h_2 < \dots < h_N = H.$$

- On each interval $I_i := (h_i, h_{i+1})$, we approximate $h \mapsto c_{\Omega_h}$ and $h \mapsto u_{\Omega_h}^c$ by **constant quantities**:

$$c_{\Omega_h} \approx c_{\Omega_{h_{i+1}}}, \text{ and } u_{\Omega_h}^c \approx u_{\Omega_{h_{i+1}}}^c \text{ on } \Omega_h, \text{ for } h \in (h_i, h_{i+1}).$$

- These approximations are used in the formulas for $P_{\text{sw}}(\Omega)$ and $P'_{\text{sw}}(\Omega)(\theta)$ to obtain the reconstructions P_N^0 and \mathcal{D}_N^0 .
- This method is **costly** in practice, since it requires a fine discretization $\{h_i\}$ of $(0, H)$ to be accurate enough.

A first-order interpolation method (I)

- On each interval I_i , we approximate $h \mapsto c_{\Omega_h}$ by a **cubic spline** $h \mapsto \tilde{c}_i(h)$ by using the data:

$$\tilde{c}_i(h_i) = c_{\Omega_{h_i}}, \quad \tilde{c}_i(h_{i+1}) = c_{\Omega_{h_{i+1}}}, \quad \tilde{c}_i'(h_i) = \left. \frac{d}{dh}(c_{\Omega_h}) \right|_{h_i},$$

$$\text{and } \tilde{c}_i'(h_{i+1}) = \left. \frac{d}{dh}(c_{\Omega_h}) \right|_{h_{i+1}}.$$

- On each interval I_i , we approximate $h \mapsto u_{\Omega_h}^c$ by:

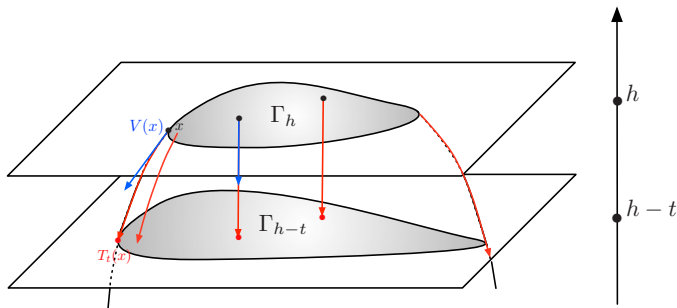
$$u_{\Omega_h}^c \approx u_{\Omega_{h_{i+1}}}^c(x) + (h_{i+1} - h) U_{\Omega_{h_{i+1}}}(x), \quad \text{a.e. } x \in \Omega_h,$$

where U_{Ω_h} is the 'derivative' of the mapping $h \mapsto u_{\Omega_h}^c \dots$ in a **suitable sense**.

- These quantities are used in the definitions of $P_{\text{sw}}(\Omega)$ and $P'_{\text{sw}}(\Omega)(\theta)$ to obtain the approximations P_N^1 and \mathcal{D}_N^1 .

A first-order interpolation method (II)

The derivative of $h \mapsto u_{\Omega_h}^c$ may be thought of as the **shape** (or **Eulerian**) derivative of the mapping $t \mapsto u_{T_t(\Omega_h)}^c$, where $T_t : \Omega_h \rightarrow \Omega_{h-t}$ is 'any' diffeomorphism from Ω_h onto Ω_{h-t} .

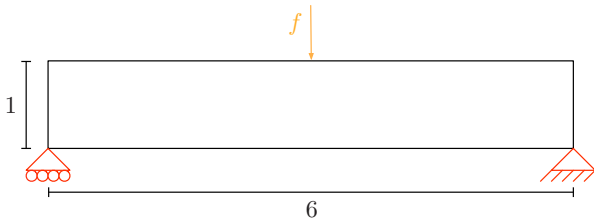


One diffeomorphism T_t mapping Ω_h onto Ω_{h-t} .

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Mechanical approach: the manufacturing compliance (I)

Still in the setting of the two-dimensional **MBB Beam** example,

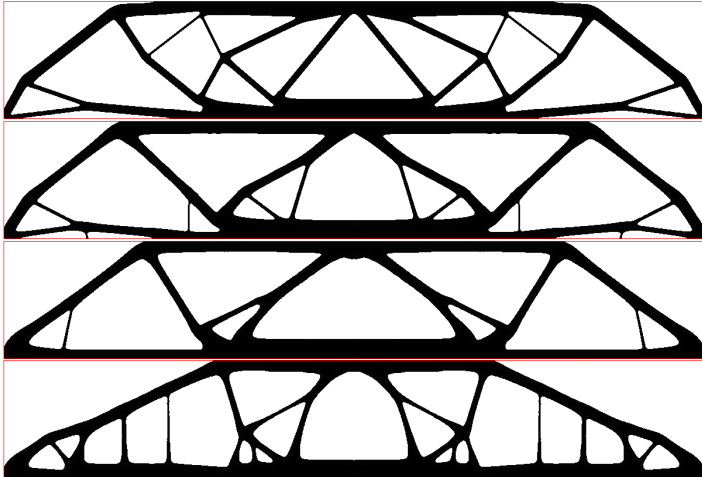


we now solve the **constrained optimization problem**:

$$\begin{aligned} \min_{\Omega} \quad & J(\Omega) \\ \text{s.t.} \quad & \text{Vol}(\Omega) \leq \alpha_v \text{Vol}(D), \\ & P_{\text{sw}}(\Omega) \leq \alpha_c P_{\text{sw}}(\Omega^*), \end{aligned}$$

where $\alpha_c \in [0, 1]$ is a user-defined tolerance, and Ω^* is the optimized shape for the compliance under volume constraint (without additive manufacturing constraint).

Mechanical approach: the manufacturing compliance (II)

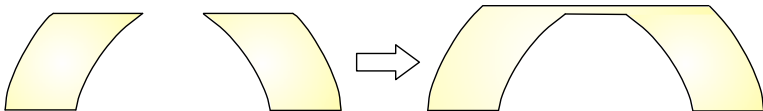


Optimized shapes for the two-dimensional MBB Beam example; (top) optimized shape Ω^ , without additive manufacturing constraints, and optimized shapes using parameters (from top to bottom) $\alpha_c = 0.50$, $\alpha_c = 0.30$, and $\alpha_c = 0.10$.*

Mechanical approach: the manufacturing compliance (III)

This new approach yields better results; yet, it raises two issues:

1. $P_{\text{sw}}(\Omega)$ inherently favors structures whose lower part is stronger.
2. The optimized shapes still show large, completely horizontal overhangs. This is a flaw in the modelling of $P_{\text{sw}}(\Omega)$, which **assumes that each layer of material is assembled instantaneously**.



*Completely flat overhangs are not so weak because of the **instantaneous layer deposition** assumption.*

Mechanical approach: the modified manufacturing compliance (I)

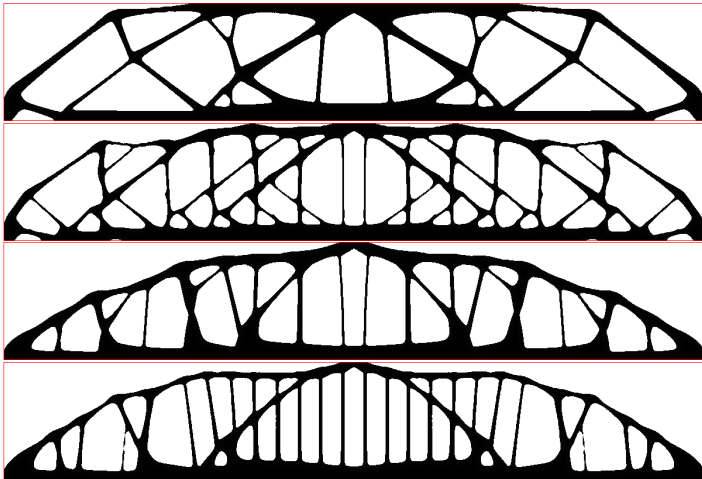
We now solve:

$$\begin{aligned} \min_{\Omega} \quad & J(\Omega) \\ \text{s.t.} \quad & \text{Vol}(\Omega) \leq \alpha_v \text{Vol}(D), \\ & P_{\text{uw}}(\Omega) \leq \alpha_c P_{\text{uw}}(\Omega^*), \end{aligned}$$

where the **modified (upper weight) manufacturing compliance** $P_{\text{uw}}(\Omega)$ brings into plays elastic displacements of the intermediate shapes $u_{\Omega_h}^a$ involving an **artificial load** concentrated on their upper side:

$$\left\{ \begin{array}{ll} -\text{div}(Ae(u_{\Omega_h}^a)) = g_h & \text{in } \Omega_h, \\ u_{\Omega_h}^a = 0 & \text{on } \Gamma_0, \\ Ae(u_{\Omega_h}^a)n = 0 & \text{on } \Gamma_h^l, \\ Ae(u_{\Omega_h}^a)n = 0 & \text{on } \Gamma_h^u, \end{array} \right. \quad \text{where } g_h(x) = \begin{cases} g & \text{if } x_d \in (h - \delta, h), \\ 0 & \text{otherwise.} \end{cases}$$

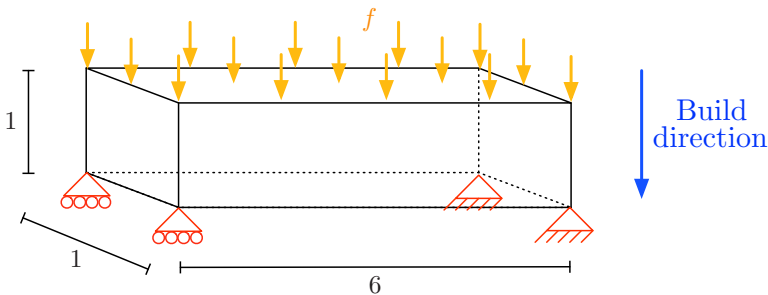
Mechanical approach: the modified manufacturing compliance (II)



Optimized 2d MBB Beams obtained using the modified manufacturing compliance $P_{af}(\Omega)$ and parameters (from top to bottom) $\alpha_c = 0.30$, $\alpha_c = 0.10$, $\alpha_c = 0.05$, and $\alpha_c = 0.03$.

Mechanical approach: the modified manufacturing compliance (III)

We now consider the design of a three-dimensional bridge.

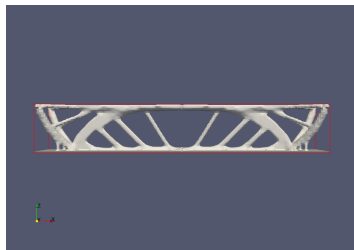
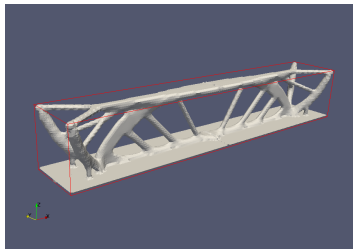
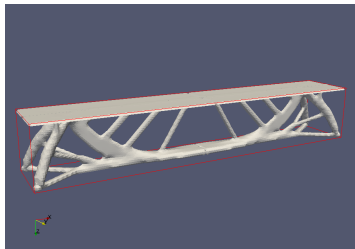


We solve the following shape optimization problem:

$$\begin{aligned} \min_{\Omega} \quad & \text{Vol}(\Omega), \\ \text{s.t.} \quad & J(\Omega) \leq J(\Omega^*), \\ & P_{\text{uw}}(\Omega) \leq \alpha_c P_{\text{uw}}(\Omega^*). \end{aligned}$$

Mechanical approach: the modified manufacturing compliance (IV)

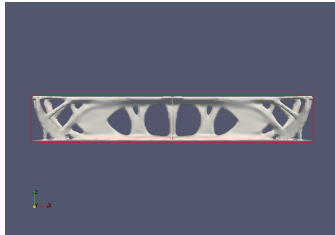
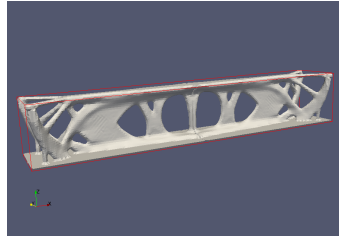
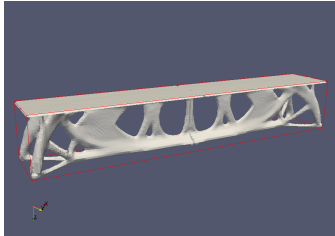
The optimized shape Ω^* without manufacturing shows several large overhangs.



Different views of the unconstrained optimized shape Ω^ .*

Mechanical approach: the modified manufacturing compliance (V)

Several small overhangs remain on the upper part of the optimized shape with the imposed manufacturing constraint $P_{uw}(\Omega)$.



Different views of the optimized shape for $\alpha_c = 0.1$.

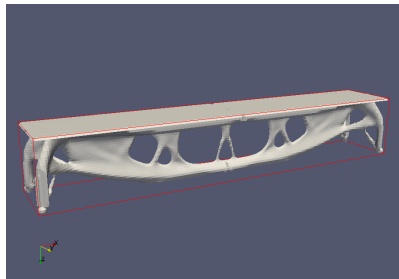
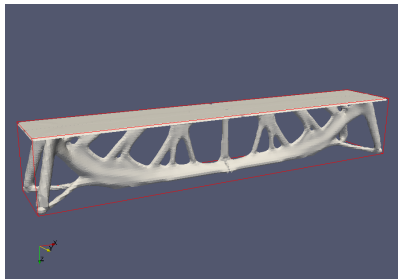
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Mixing geometric and mechanical formulations (I)

To remedy these completely flat regions caused by the **instantaneous layer deposition**, we consider the following optimization problem, which mixes **geometric** and **mechanical** constraints:

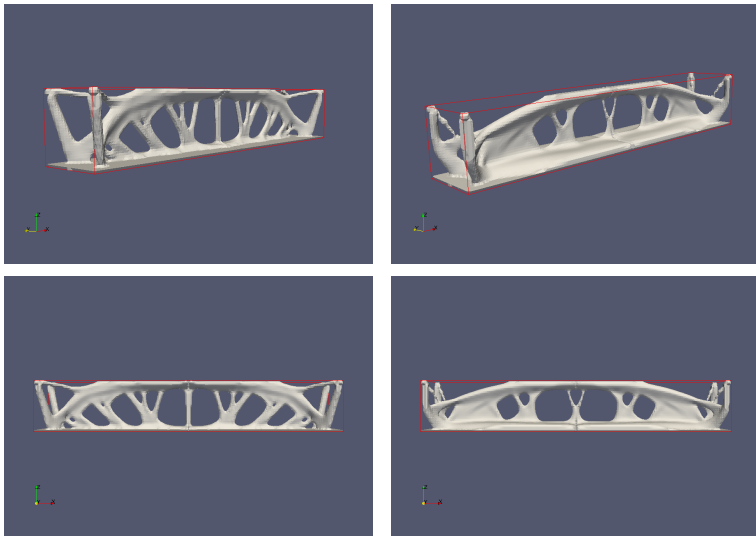
$$\begin{aligned} \min_{\Omega} \quad & (1 - \alpha_g) \frac{\text{Vol}(\Omega)}{\text{Vol}(\Omega^*)} + \alpha_g \frac{P_g(\Omega)}{P_g(\Omega^*)}, \\ \text{s.t.} \quad & J(\Omega) \leq J(\Omega^*). \\ & P_{\text{uw}}(\Omega) \leq \alpha_c P_{\text{uw}}(\Omega^*). \end{aligned}$$

Mixing geometric and mechanical formulations (II)



Optimized 3d bridges for the combined geometrically and mechanically constrained problem, setting (left) $\alpha_c = 0.10$, $\alpha_g = 0.10$, and (right) $\alpha_c = 0.10$, $\alpha_g = 0.90$.

Mixing geometric and mechanical formulations (III)








Different views of the optimized shapes with (left) $\alpha_c = 0.10$, $\alpha_g = 0.10$, and (right) $\alpha_c = 0.10$, $\alpha_g = 0.90$.

Thank you !

Thank you for your attention!

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